## A Level Physics Online

## OCR B Physics - H557

## Module 5: Rise and Fall of the Clockwork Universe



| You should be able to demonstrate and show your understanding of: | Progress and understanding: |  |  |  |
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| We can derive an equation to relate linear and angular velocity. If a particle is moving in a circular orbit with linear (tangential) velocity, $v$, it moves a distance $\mathrm{I}=\mathrm{r} \theta$ (from above) in a time, t ; $v=\frac{d l}{d t}=\frac{d(r \theta)}{d t}=r \frac{d \theta}{d t}=r \omega \quad \text { so, } \quad v=\omega r$ <br> Note: $r$ is taken outside the differential because it is a constant multiplier |  |  |  |  |
| Using this new equation, we can substitute it into the above equations for centripetal force and acceleration to get; $F=m r \omega^{2} \quad a=\omega^{2} r$ |  |  |  |  |
| The gravitational force is given by the equation below, where $G$ is the gravitational constant and $\mathrm{G}=6.67 \times 10^{-11} \mathrm{Nm}^{2} \mathrm{~kg}$. There is a negative sign because the force due to gravity is always attractive $F_{g r a v}=-\frac{G M m}{r^{2}}$ <br> Where $M$ and $m$ are the masses concerned in orbit and $r$ is the distance between the mass concerned and the centre of orbit |  |  |  |  |
| Test mass: A mass small enough so that it does not affect the surrounding gravitational field |  |  |  |  |
| If an object is put into a gravitational field say it is subject to a force, NOT it feels a force |  |  |  |  |
| Gravitational Field Strength, g (Units: $\mathrm{Nkg}^{-1}$ ): The magnitude and direction of the force on 1 kg at a given point in a gravitational field. Field lines represent the direction and magnitude of it on a diagram; $g=-\frac{G M}{r^{2}}$ <br> Note: $g$ varies with an inverse square law as shown by the equation |  |  |  |  |
| Gravitational Field inside the Earth: Proportional to the distance from the centre. The mass of Earth above a point inside the planet does not have a gravitational effect. Considering a point, P, inside the Earth; $g_{P}=-\frac{G M_{p}}{r_{P}^{2}}$ <br> where $M_{P}=\rho \frac{4}{3} \pi r_{P}^{3}$ is the mass of Earth below $P$, hence $g_{P}=-\rho \frac{4}{3} G \pi r_{P}$ |  |  |  |  |
| Uniform Field: A field where there is the same magnitude and direction of field everywhere. For example, near the Earth's surface is a uniform field |  |  |  |  |


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| Potential in a Uniform Field: For a falling particle; $\Delta E_{P}=\Delta E_{k} \quad \rightarrow \quad \Delta m g h=\Delta \frac{1}{2} m v^{2} \quad \rightarrow \quad \Delta g h=\Delta \frac{1}{2} v^{2}$ |  |  |  |  |
| Geostationary Satellites: Stay in the same place in the sky. Orbit the Earth above the equator once a day. An orbit that is too low will result in it moving too fast, an orbit that is too high will result in it moving too slow. A stable orbit will be achieved if the velocity of the satellite equals the rate of Earth's rotation, this occurs if the centripetal force equals the gravitational force; $\begin{gathered} \frac{m v^{2}}{r}=\frac{G M m}{r^{2}} \text { then substitute in } v=\frac{2 \pi r}{T} \\ \frac{4 \pi^{2} r^{2}}{T^{2}}=\frac{G M}{r} \rightarrow r^{3}=\frac{G M T^{2}}{4 \pi^{2}} \text { this results in; } \\ r^{3} \propto T^{2} \end{gathered}$ <br> This is Kepler's third law |  |  |  |  |
| The orbital period of the geostationary satellite can also be determined; $\begin{gathered} T=\frac{2 \pi r}{v} \rightarrow T^{2}=\frac{4 \pi^{2} r^{2}}{v^{2}} \text { then substitute in } v^{2}=\frac{G M}{r} \\ T^{2}=\frac{4 \pi^{2} r^{3}}{G M} \text { then substitute in } g=\frac{G M}{r^{2}} \\ T=2 \pi \sqrt{\frac{r}{g}} \end{gathered}$ |  |  |  |  |
| Gravitational Potential, V (Units: $\mathrm{Jkg}^{-1}$ ): The gravitational energy per kg of material. For a uniform near Earth, $\Delta V=9.81 \mathrm{Jkg}^{-1}$ for every metre raised $V=-\frac{G M}{r}$ |  |  |  |  |
| Equipotential Surfaces: A continuous surface of the same gravitational potential. In a uniform field, the spacing between equipotentials is equal. Along an equipotential, not gravitational force acts because the potential energy is not changing hence no work is done. The direction of the gravitational field is always perpendicular to the equipotentials. The change in potential between two points (on different equipotentials) is the same irrespective of the route taken. This is true for uniform and non-uniform fields |  |  |  |  |
| Gravitational Potential Wells: For calculations concerning celestial objects, zero is set as the potential energy per kg at infinity. Gravitational potential increases the further you get from Earth. A more massive object has a more negative potential at its surface. For a spacecraft travelling from a point of net gravitational neutrality to Earth, its kinetic energy increases and its potential energy decreases as it 'falls down' the potential well. A point of net gravitational neutrality is where there is zero net force acting |  |  |  |  |


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| Escape Velocity: <br> For a particle to escape a gravitational potential well, $E_{T O T}=E_{P}+E_{k} \geq 0 \quad \rightarrow \quad \frac{1}{2} m v^{2}=\frac{G M m}{r^{2}}$ <br> Rearrange for $v$ gives the minimum escape velocity, $v_{e s c}=\sqrt{\frac{2 G M}{r}}=\sqrt{2} v_{o r b}$ <br> To aid in escaping a potential well there are ways of gaining energy; the Earth is rotating so the launch site has kinetic energy already. A gravity assist/slingshot is where a body gains kinetic energy through interacting with another planet |  |  |  |  |
| Orbital Velocity: If a body has a higher velocity than the orbital velocity, it will move to a higher orbit. If the orbital velocity is greater than the escape velocity for the radius of orbit, the body will leave the planet's orbit $v_{o r b}=\sqrt{\frac{G M}{r}}$ |  |  |  |  |
| Radial Field Equations: Some of these have already been covered above, ut just to make explicit that these apply to radial fields (where the field lines meet at a central point), they are copied here; $F_{\text {grav }}=-\frac{G M m}{r^{2}} \quad g=-\frac{G M}{r^{2}} \quad E_{P}=-\frac{G M m}{r} \quad V=-\frac{G M}{r}$ <br> The only one not covered previously is the equation for gravitational potential energy. A simple way to think about it is to use the work equation, $W=F r$, where $F$ is the gravitational force, so an $r$ on the bottom cancels with the one on the top. <br> For a radial field, the spacing of equipotentials separated by equal potential, V , increases with distance from a point mass. Greater spacing means a weaker field |  |  |  |  |
| Potential Gradient: If a mass moves a distance $\Delta r$ (not along an equipotential), its potential changes by $\Delta \mathrm{V}$. The force required to move the mass vertically upwards at a constant velocity is; $\begin{aligned} & F=-m g \text { and the work done is, } W=F d=\Delta E_{P}=V m \text { so; } \\ & -\quad-m g \Delta r=\Delta V m ; \\ & \qquad \begin{array}{c} \Delta=-\frac{\Delta V}{\Delta r} \text { and as } \Delta r \rightarrow 0 ; \\ g=\frac{d V}{d r} \end{array} \end{aligned}$ <br> This is the potential gradient (for a radial field). The radial component of the field at any point in space equals the negative of the potential gradient at that point |  |  |  |  |
| A g-r graph is in the fourth quadrant as $g$ is always negative and $r$ is positive. The area of the graph between the graph and the x -axis gives $\Delta \mathrm{V}$ |  |  |  |  |


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| An F-r graph is also in the fourth quadrant, for similar reasons. The area of <br> the graph between the graph and the x-axis gives $\Delta E_{P}$ |  |  |  |  |
| A V-r graph is in the fourth quadrant, for similar reasons. The gradient of the <br> graph gives -g (The gradient is positive so to find g it is the negative of the <br> gradient) |  |  |  |  |
| An Ep-r graph is in the fourth quadrant, for similar reasons. The gradient of <br> the graph gives -F (The gradient is positive so to find F it is the negative of <br> the gradient) |  |  |  |  |

